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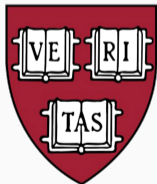
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Deconfined quantum criticality and a gapless \mathbb{Z}_2 spin liquid in the square lattice antiferromagnet

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Numerical evidence for Néel/SL/VBS transition in J_1/J_2 model

$$H = J_1 \sum_{\langle i,j \rangle} \mathbf{S}_i \cdot \mathbf{S}_j + J_2 \sum_{\langle\langle i,j \rangle\rangle} \mathbf{S}_i \cdot \mathbf{S}_j$$

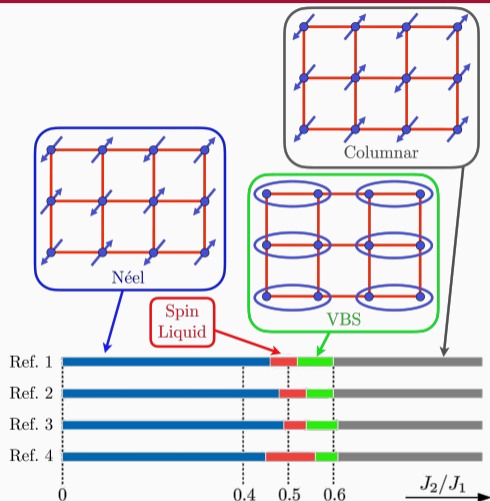
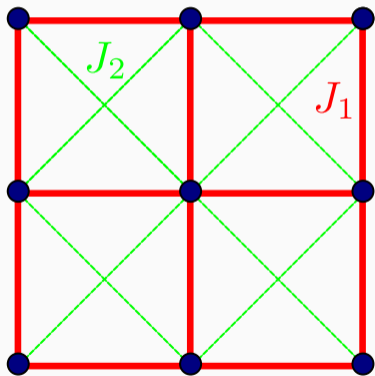
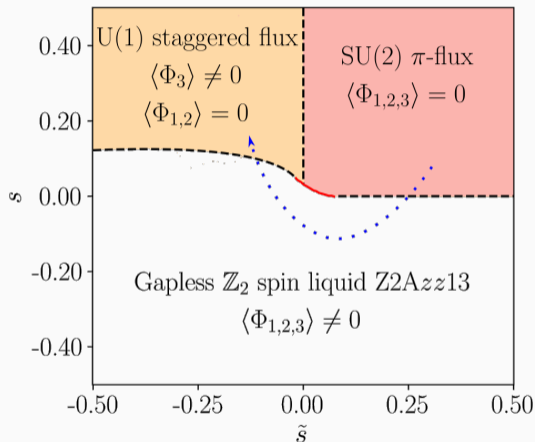


Figure adapted from Becca 2020

L. Wang and Sandvik 2018; Ferrari and Becca 2020; Nomura and Imada 2020; Liu et al. 2020.

Deconfined critical theory for Néel/ \mathbb{Z}_2 and \mathbb{Z}_2 /VBS transitions

- Transition described by Higgs condensation breaking $SU(2) \rightarrow \mathbb{Z}_2$ or $U(1) \rightarrow \mathbb{Z}_2$
- Both $SU(2)$ (C. Wang et al. 2017) and $U(1)$ (Song et al. 2019) spin liquids unstable to Néel or VBS order on square lattice



Large N_f critical theory of $U(1)$ Higgs transition

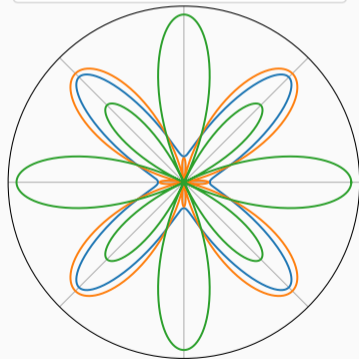
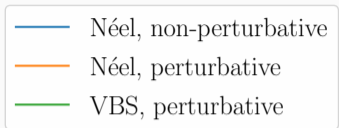
$$\mathcal{L} = \mathcal{L}_{N=4 \text{ QED}_3} + \mathcal{L}_{\text{Yukawa}} + \overbrace{\Phi \bar{\psi} \mu^y (\gamma^x D_x - \gamma^y D_y) \psi}^{\text{Velocity anisotropy}}$$

Large N_f critical theory of $U(1)$ Higgs transition

$$\mathcal{L} = \mathcal{L}_{N=4 \text{ QED}_3} + \mathcal{L}_{\text{Yukawa}} + \overbrace{\Phi \bar{\psi} \mu^y (\gamma^x D_x - \gamma^y D_y) \psi}^{\text{Velocity anisotropy}}$$

Fixed point to $\mathcal{O}(N_f^{-1})$:

- $\Phi_c \approx 0.46$
- $z \approx 1 + 0.23 N_f^{-1}$
- $\eta_{\text{Néel}} = 0.065 N_f^{-1}$
- $\eta_{\text{VBS}} = -0.01 N_f^{-1}$



Large N_f critical theory of SU(2) Higgs transition

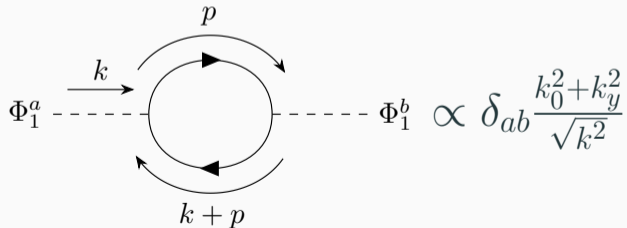
$$\mathcal{L} = \mathcal{L}_{N=2 \text{ QCD}_3} + s (\Phi_1^a \Phi_1^a + \Phi_2^a \Phi_2^a) + \Phi_1^a \bar{\psi} \gamma^x \mu^z \sigma^a \psi + \Phi_2^a \bar{\psi} \gamma^y \mu^x \sigma^a \psi$$

Large N_f critical theory of $SU(2)$ Higgs transition

$\mathcal{L} = \mathcal{L}_{N=2 \text{ QCD}_3} + s (\Phi_1^a \Phi_1^a + \Phi_2^a \Phi_2^a) + \Phi_1^a \bar{\psi} \gamma^x \mu^z \sigma^a \psi + \Phi_2^a \bar{\psi} \gamma^y \mu^x \sigma^a \psi$
 $\bar{\psi} \not{\partial} \psi + \Phi_1^a \bar{\psi} \gamma^x \mu^z \sigma^a \psi$ invariant under subsystem "symmetry":

$$\psi \rightarrow e^{i f_a(x) \mu^z \sigma^a} \psi$$

$$\Phi_1^a \rightarrow U_{ab}^{-1} \Phi_1^b + \partial_x f_a(x)$$



Subsystem symmetries control universal behavior of critical theory

Must keep subleading “dangerously irrelevant” terms, $K (\Phi_1^a \partial^2 \Phi_1^a + \Phi_2^a \partial^2 \Phi_2^a)$, in order to regulate our theory.



$$\Sigma(k) \approx -\frac{3}{\pi^2 N_f} k \ln^2(K|k|)$$

$$G_{\text{Néel}}(r) \sim \frac{1}{r^\alpha} \exp \left[-\frac{12}{\pi^2 N_f} \ln^2(r) \right]$$

$$G_{\text{VBS}}(r) \sim \frac{1}{r^\beta} \exp \left[-\frac{6}{\pi^2 N_f} \ln^2(r) \right]$$

Conclusions and future directions

- Deconfined critical theories between spin liquids and Néel/VBS order captured by Higgs transitions to proximate unstable SLs
- Critical theories exhibit Lorentz symmetry breaking, violation of standard scale invariance, Néel/VBS asymmetry

Open questions:

- Numerical signatures?
- Other models with emergent subsystem symmetries at criticality?

