

# A sign-problem-free effective model of triangular lattice antiferromagnetism

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Henry Shackleton

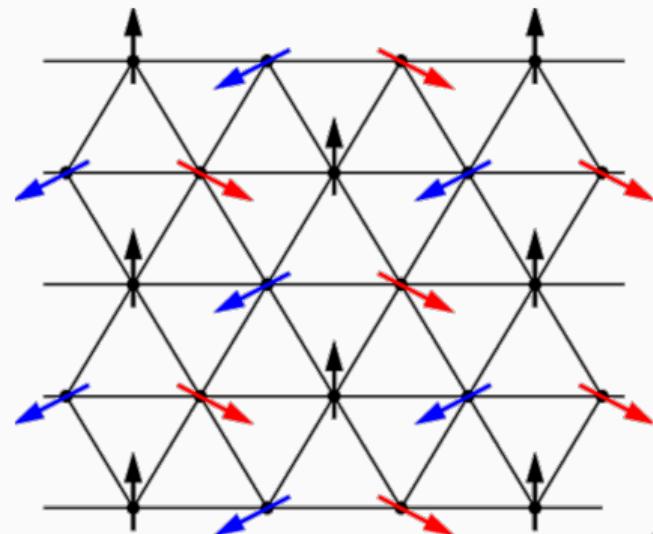
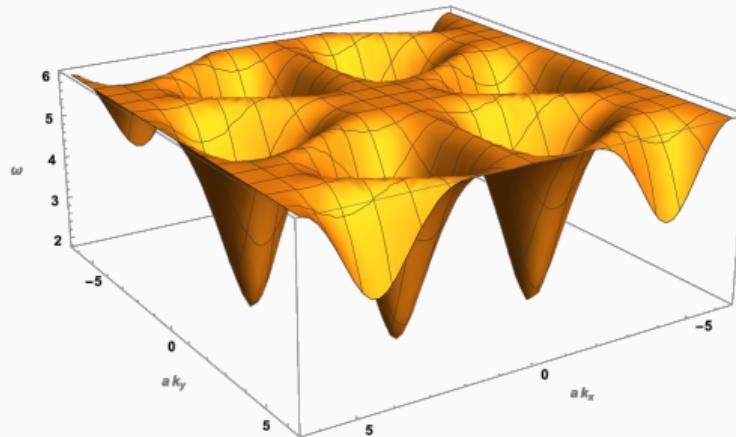
March 6, 2023

Harvard University



# Frustrated antiferromagnetism leads to exotic quantum phases

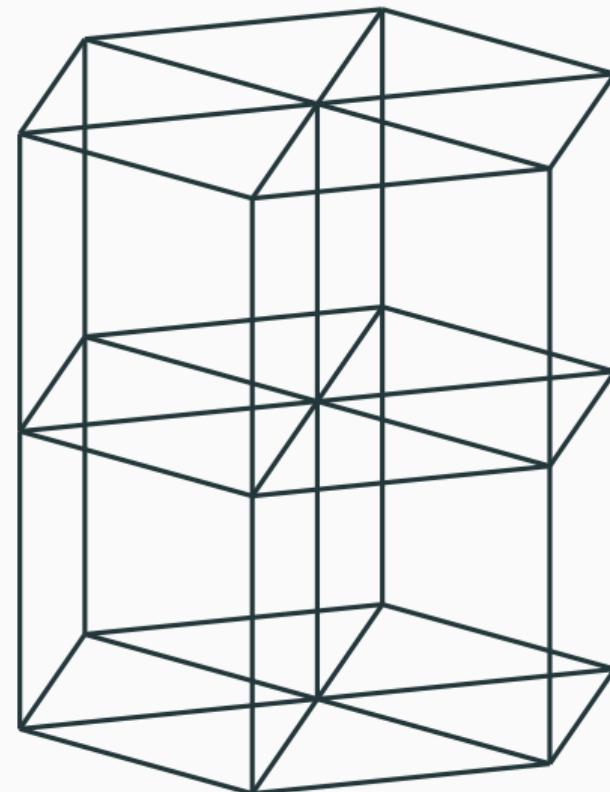
- $H = \sum_{ij} J_{ij} \mathbf{S}_i \cdot \mathbf{S}_j$ , trivial phase for half-integer spin precluded by LSM theorem
- Study *effective* model starting from bosonic spinon construction,  $\mathbf{S}_i = b_{i\alpha}^\dagger \boldsymbol{\sigma}^{\alpha\beta} b_{i\beta}$



# Effective model of triangular lattice antiferromagnetism

O(4) bosonic spinons  $z_{i\alpha}$ ,  $\alpha = 1, 2$   
coupled to an *odd*  $\mathbb{Z}_2$  gauge field  $s_{j,j+\hat{\mu}}$

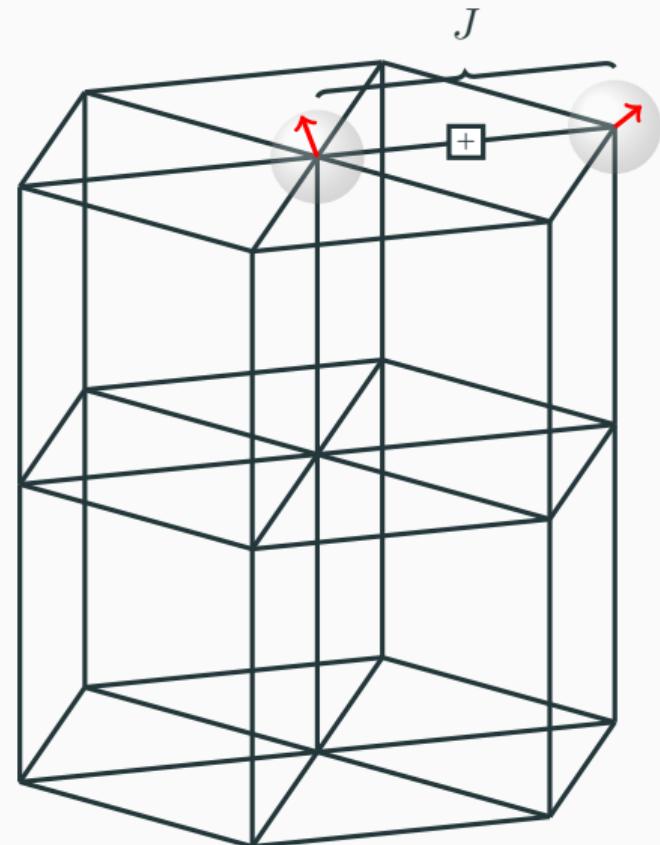
$$H[z_\alpha, s] = -\frac{J}{2} \sum_{\langle j, \mu \rangle} s_{j,j+\hat{\mu}} (z_{j,\alpha}^* z_{j+\hat{\mu},\alpha} + \text{c.c})$$
$$- K \sum_{\Delta \square} \prod_{\Delta \square} s_{j,j+\hat{\mu}} + i\pi \underbrace{\sum_j s_{j,j+\hat{\tau}}}_{\text{Berry phase}}$$



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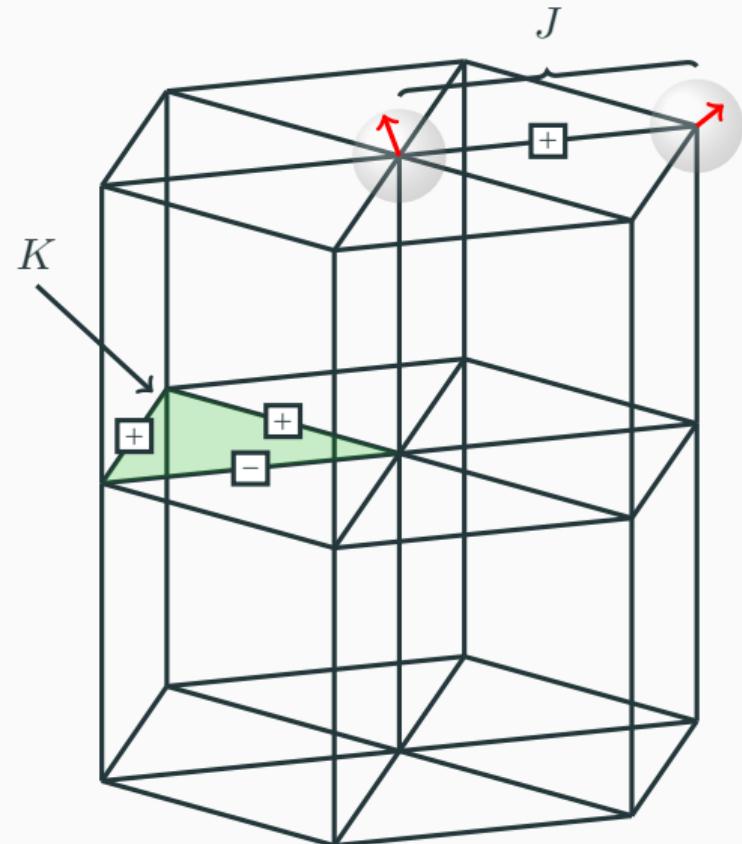
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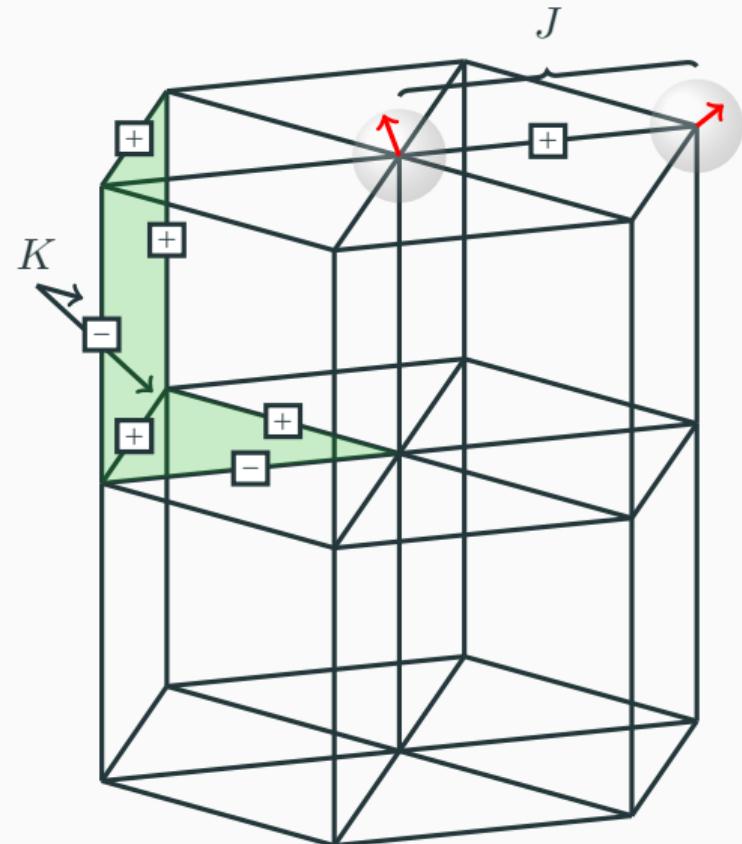
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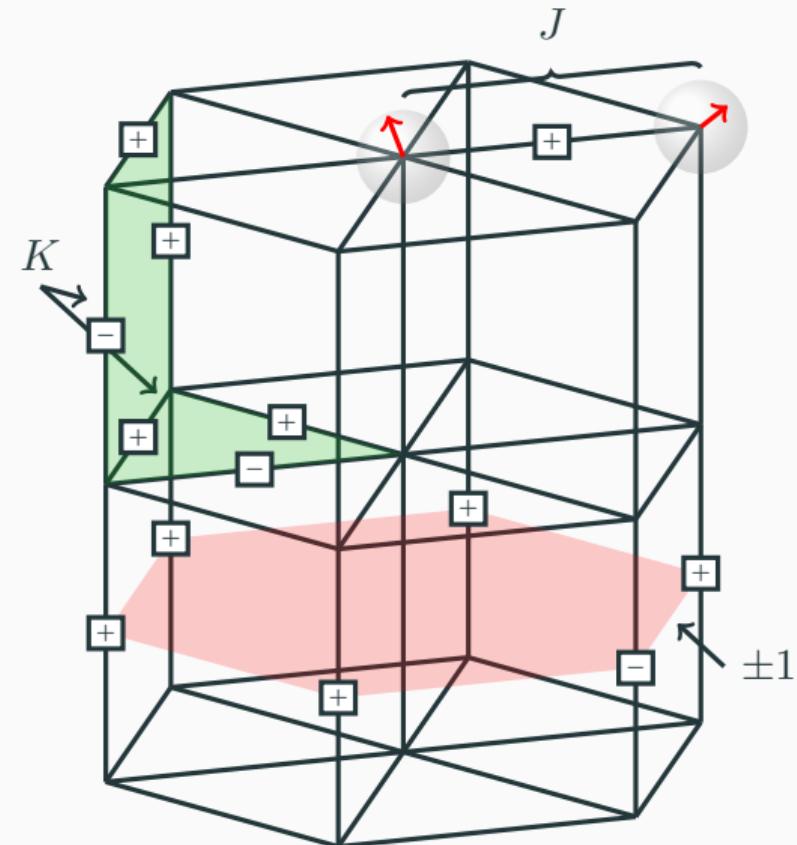
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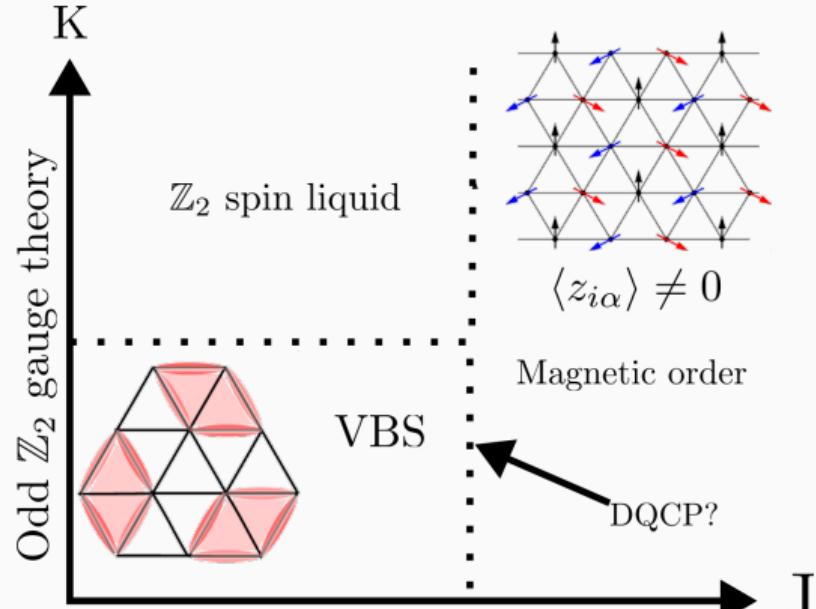
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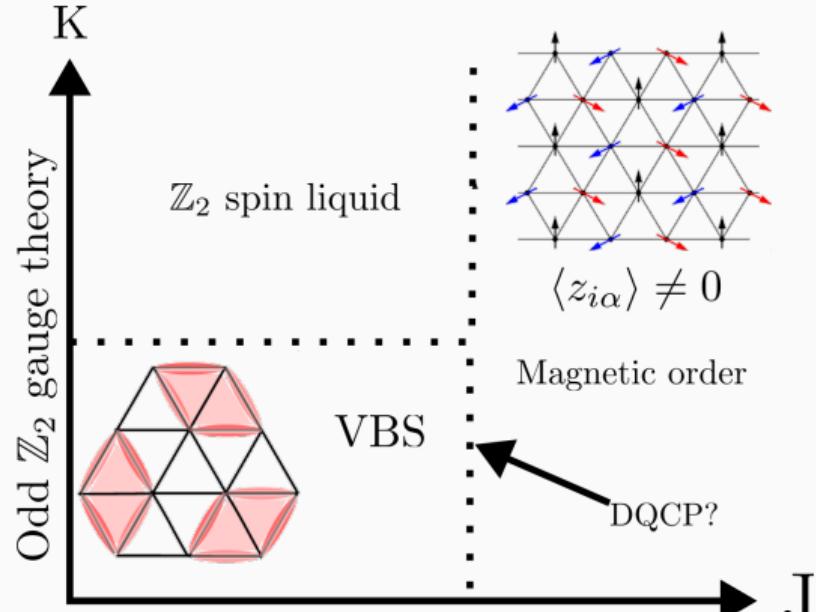
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## Session D54: Quantum Spin Liquids: Theory and Computation

3:00 PM–5:48 PM, Monday, March 6, 2023  
Room: Room 306

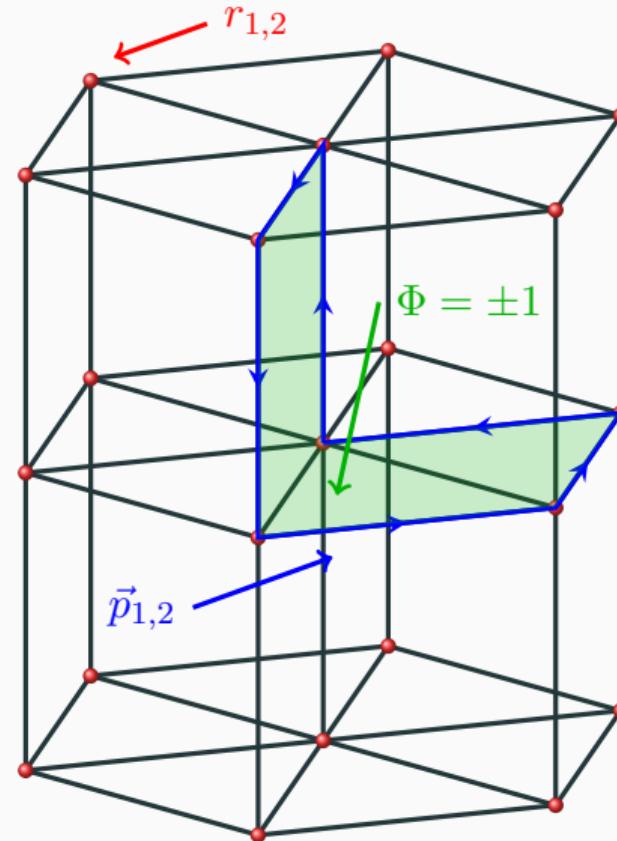
**Abstract:** D54.00007 : Quantum Electrodynamics in 2+1 Dimensions as the Organizing Principle of Triangular Lattice Antiferromagnets  
4:12 PM–4:24 PM

**Presenter:**  
Andreas M Läuchli  
(Paul Scherrer Institute)

[← Abstract →](#)

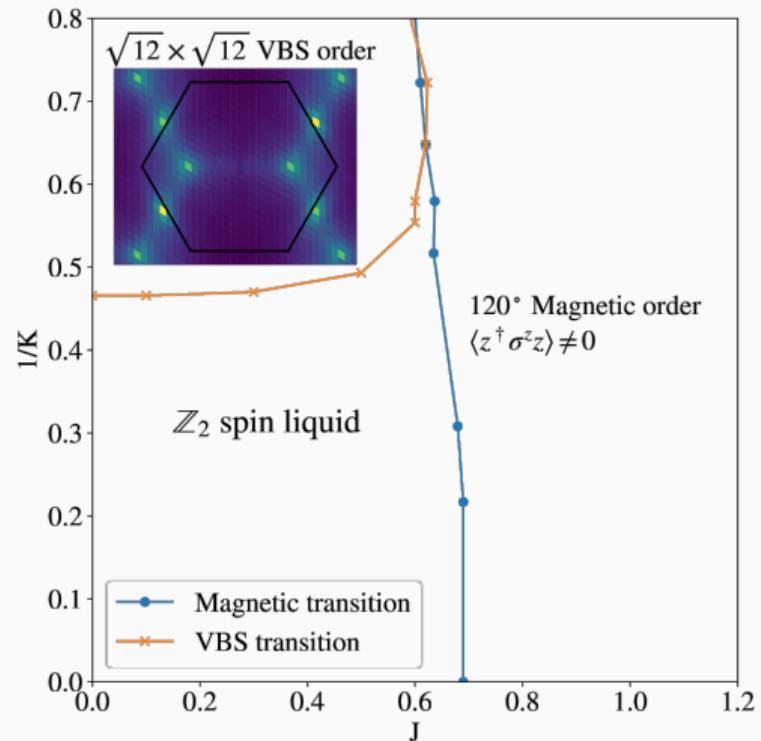
## Berry phase requires sign-problem-free mapping

- $z_{j,\alpha} = r_{j,\alpha} e^{i\theta_{j,\alpha}}$ , angular variables recast into closed current loops
- Non-trivial constraint between allowed gauge flux  $\Phi$  and current loops
- Berry phase reflected in *geometric frustration* of gauge flux



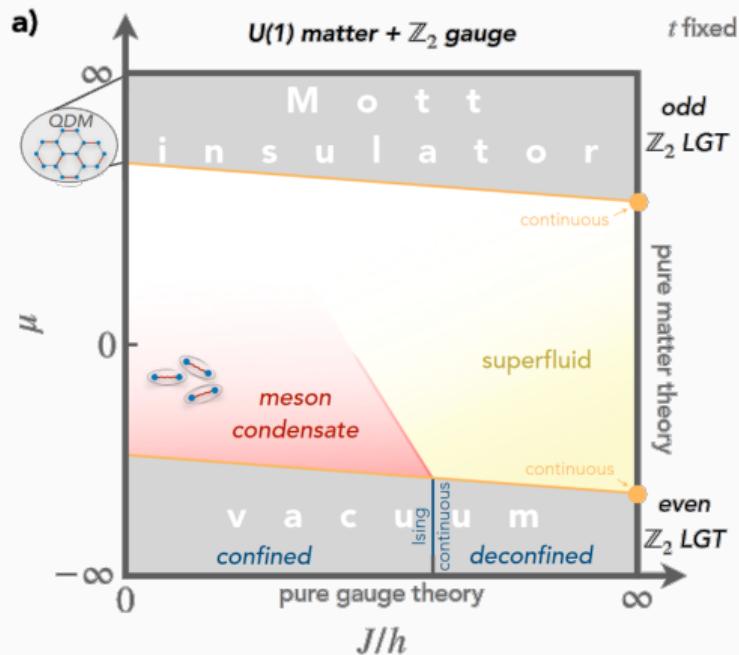
# Phase diagram captures spin liquid and ordered phases

- Preliminary data indicates direct transition between VBS and magnetic order
- System sizes restricted to multiples of 12 to admit VBS order
- Independent spinon/gauge cluster updates insufficient to resolve transition at large system sizes



## Future work

- Improved global updates for studying DQCP - can we find signatures of  $\text{QED}_3$ ?
- Generalization to continuous-time simulations
- Application of sign-problem-free mapping to more systems - Bose Hubbard +  $\mathbb{Z}_2$  at finite density



## Full sign-problem-free Hamiltonian

$$\mathcal{Z} = \sum_{h_{\bar{j},\alpha,\mu}=-\infty}^{\infty} \prod_{j\alpha} \int_0^1 r_{j,\alpha} \, dr_{j,\alpha} \delta \left( \sum_{\alpha} r_{j,\alpha}^2 - 1 \right) \exp(-H[r_{\alpha}, h_{\alpha}])$$

$$H[r_{\alpha}, h_{\alpha}] = \sum_{\langle j, \mu \rangle} \left[ -\ln I_{p_{j,\alpha,\mu}}(J r_{j,\alpha} r_{j+\hat{\mu},\alpha}) + K_d \varepsilon_{\bar{j},\mu} \sigma_{\bar{j},\mu} \right]$$

$$\sigma_{\bar{j},\mu} = 2 \sum_{\alpha} h_{\bar{j},\alpha,\mu} \bmod 2 - 1$$

$$e^{2K_d} = \tanh K = \frac{e^K - e^{-K}}{e^K + e^{-K}}$$